# $\Gamma$ -convergence of Oscillating Thin Obstacles

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#### Abstract

Consider the minimum problems of obstacle type

$$\min\left\{\int_{\Omega} |Du|^2 dx : u \ge \psi_{\varepsilon} \text{ on } P, \ u = 0 \text{ on } \partial\Omega\right\},\$$

as  $\varepsilon \to 0$ . Here  $\psi_{\varepsilon}$  is a periodic function of period  $\varepsilon$ , constructed from an appropriately rescaled fixed function and  $P \subset \subset \Omega \subset \mathbb{R}^n$  is a subset of the hyper-plane  $\{x \in \mathbb{R}^n : x \cdot \eta = 0\}$ . We assume  $n \geq 3$  and that the normal  $\eta$  satisfies a generic condition that guarantees certain ergodic properties of the quantity

$$\#\left\{k\in\mathbb{Z}^n:P\cap\{x:|x-\varepsilon k|<\varepsilon^{n/(n-1)}\}\right\}.$$

Under these hypotheses we compute explicitly the limit functional of the obstacle problem above, which is of the type

$$H^1_0(\Omega) \ni u \mapsto \int_{\Omega} |Du|^2 dx + \int_P G(u) d\sigma.$$

# 1 Preliminaries and Main Result

### 1.1 Introduction of the Problem

We consider an obstacle problem in a domain  $\Omega \subset \mathbb{R}^n$  for  $n \geq 3$ . The obstacle is the restriction to a hyper-plane of a rescaled, periodically extended function. The given data in the problem is

- 1. A domain  $\Omega$  in  $\mathbb{R}^n$ ,  $n \geq 3$ , i.e. a bounded, open, connected subset of  $\mathbb{R}^n$ .
- 2. A continuous function  $\psi$  with compact support in  $B_{1/2} = \{x \in \mathbb{R}^n : |x| < 1/2\}.$
- 3. A hyper-plane  $\Pi = \{x \in \mathbb{R}^n : x \cdot \eta = 0\}$  with unit normal  $\eta = (\eta_1, \ldots, \eta_n)$  such that  $e_n \notin \Pi \iff \eta_n \neq 0$ .

Note that for any  $E \subset \mathbb{R}^n$ ,  $P := E \cap \Pi$  can be represented as

$$P = \{ (x', \alpha x') : x' \in H \},$$
(1)

where  $x' = (x_1, \dots, x_{n-1}), x = (x', x_n),$ 

$$H = \operatorname{proj}_{\mathbb{R}^{n-1}} P$$

and

$$\alpha = (\alpha_1, \dots, \alpha_{n-1}), \quad \alpha_i = \frac{-\eta_i}{\eta_n}.$$

Let  $Q_{\varepsilon} = (-\varepsilon/2, \varepsilon/2)$ , and for any  $k \in \mathbb{Z}^n$ , let  $Q_{\varepsilon}^k = Q_{\varepsilon} + \varepsilon k$ . Similarly,  $B_{r_{\varepsilon}}^k$  denotes the ball of radius  $r_{\varepsilon}$  and center  $\varepsilon k$ , i.e.  $B_{r_{\varepsilon}}^k = B_{r_{\varepsilon}} + \varepsilon k$ . From  $\psi$  we construct the oscillating function  $\psi_{\varepsilon}$ , given by

$$\psi_{\varepsilon}(x) = \begin{cases} \psi(a_{\varepsilon}^{-1}(x - \varepsilon k)), & \text{if } x \in Q_{\varepsilon}^{k} \cap \Pi, \\ -\infty, & \text{otherwise,} \end{cases}$$
(2)

where

$$a_{\varepsilon} = \varepsilon^{n/(n-1)}.$$
(3)

**Remark 1.** From the definition of  $\psi_{\varepsilon}$  it can be seen that  $\psi_{\varepsilon}(x) > -\infty$  if and only if

$$x \in \{a_{\varepsilon}\{y: \psi(y) > -\infty\} + \varepsilon k\} \cap \Pi, \text{ for some } k \in \mathbb{Z}^n.$$

For this reason it needs to be determined how often  $\Pi$  intersects a neighbourhood of size comparable to  $a_{\varepsilon}$  of the lattice points  $\{\varepsilon k\}_{k\in\mathbb{Z}^n}$ . This is possible in  $n \geq 3$  dimensions, using the theory of uniform distribution of sequences. In general, this is possible when  $a_{\varepsilon}$  is not "too small". When n = 2 we would have to choose a much smaller  $a_{\varepsilon}$ , due to the logarithmic nature of the fundamental solution of the laplacian. For this reason we cannot include the two dimensional case.

For any Borel subset  $\mathcal{B}$  of  $\Omega$  and  $u \in H_0^1(\Omega)$ , set

$$F_{\psi_{\varepsilon}}(u, \mathcal{B}) = \begin{cases} 0, & \text{if } u \ge \psi_{\varepsilon} \text{ q.e. on } \mathcal{B}, \\ \infty, & \text{otherwise,} \end{cases}$$
(4)

where q.e. is short for quasi everywhere, i.e. everywhere except for a set of zero capacity. Note that  $\mathcal{B} \mapsto F_{\psi_{\varepsilon}}(u, \mathcal{B})$  only depends on  $\mathcal{B} \cap \Pi$ . Our main goal is to determine the asymptotic behaviour, as  $\varepsilon \to 0$ , of minimizers of the functional

$$J_{\varepsilon}(u) = \int_{\Omega} |Du|^2 dx + F_{\psi_{\varepsilon}}(u, \mathcal{B}).$$
(5)

#### **1.2** The Notion of $\Gamma$ -convergence

**Definition 1** ( $\Gamma$ -convergence). A sequence of functionals  $J_{\varepsilon}$  on a topological space V is said to  $\Gamma$ -converge to the functional  $J_0$  if the following hold for all  $v \in V$ :

(i) whenever  $v_{\varepsilon} \to v$  in V,

$$J_0(v) \le \liminf_{\varepsilon \to 0} J_{\varepsilon}(v_{\varepsilon}),$$

(ii) there exists a sequence  $\{v_{\varepsilon}\}_{\varepsilon}$  such that  $v_{\varepsilon} \to v$  in V and

$$J_0(v) \ge \limsup_{\varepsilon \to 0} J_{\varepsilon}(v_{\varepsilon}).$$

The functional  $J_0$  is called the  $\Gamma$ -limit of  $J_{\varepsilon}$ .

**Remark 2.** It follows easily from this definition that if  $J_{\varepsilon} \Gamma$ -converges to  $J_0$ , if  $v_{\varepsilon} \in V$  solves  $\inf_V J_{\varepsilon}(v) = J_{\varepsilon}(v_{\varepsilon})$  and if  $v_{\varepsilon} \to v_0$  in V, then  $J_0(v_0) = \inf_V J_0(v)$ . Indeed,  $J_0(v_0) \leq \liminf_{\varepsilon \to 0} J_{\varepsilon}(v_{\varepsilon})$  by (i), and for any other  $v \in V$ , there exists according to (ii) a sequence  $\{\bar{v}_{\varepsilon}\}_{\varepsilon}$  converging to v in V such that  $J_0(v) \geq \limsup_{\varepsilon \to 0} J_{\varepsilon}(\bar{v}_{\varepsilon})$ . Since  $J_{\varepsilon}(v_{\varepsilon}) \leq J_{\varepsilon}(\bar{v}_{\varepsilon})$ ,  $J_0(v_0) \leq \liminf_{\varepsilon \to 0} J_{\varepsilon}(v_{\varepsilon}) \leq J_0(v)$ , which proves the claim.

Next we quote a theorem of De Giorgi, Dal Maso and Longo from [4]. It is a compactness result for quadratic functionals of obstacle type and states that there is a representation theorem for the  $\Gamma$ -limits of these functionals. The compactness part of the theorem is valid for obstacle functionals for which there exists a sequence  $u_{\varepsilon} \in H_0^1(\Omega)$  such that both  $J_{\varepsilon}(u_{\varepsilon})$  and  $||u_{\varepsilon}||_{H_0^1(\Omega)}$  are bounded. This will be true if we assume that the set  $\mathcal{B}$  in (4) is compactly contained in  $\Omega$ . For the formulation below we refer to Attouch and Picard [1].

**Theorem 1** ([4]). There is a rich family  $\mathcal{R}$  of Borel subsets of  $\Omega$  such that for every  $\mathcal{B} \in \mathcal{R}$  satisfying  $\mathcal{B} \subset \subset \Omega$ , the sequence of functionals

$$J_{\varepsilon}(u) = \int_{\Omega} |Du|^2 dx + F_{\psi_{\varepsilon}}(u, \mathcal{B})$$
(6)

has a subsequence that  $\Gamma$ -converges to

$$J_0(u) = \int_{\Omega} |Du|^2 dx + \int_{\mathcal{B}} f(x, u) d\mu + \nu(\mathcal{B}), \tag{7}$$

where  $\mu$  and  $\nu$  are positive Radon measures,  $\mu \in H^{-1}(\Omega)$  and f(x, u) is convex and monotone non-increasing with respect to u.

**Remark 3.** It may be assumed that  $\nu = 0$ , c.f. [1], Theorem 4.1. We refer to [1] for the definition of a rich family of Borel sets. However, we would like to point out that a rich family  $\mathcal{R}$  of the Borel sets of  $\Omega$  is dense in the Borel sets, in the sense that for any Borel sets A, B such that  $\overline{A} \subset intB$ , there exists  $E \in \mathcal{R}$  such that  $\overline{A} \subset intE \subset \overline{E} \subset intB$ .

#### 1.3 Main Theorem

Next we define the functional that is the  $\Gamma$ -limit of  $J_{\varepsilon}$  in (5). For any  $\lambda \in \mathbb{R}$ , let

$$\psi^{\lambda}(x) = \begin{cases} \psi(x), & x \in \{P + \lambda\eta\}, \\ -\infty, & \text{otherwise,} \end{cases}$$
(8)

and set

$$g^{\lambda}(t) = \min\left\{\int_{\mathbb{R}^n} |Dv|^2 dx : v - t \in \mathcal{D}^{1,2}(\mathbb{R}^n), \ v \ge \psi^{\lambda} \text{ q.e. on } \mathbb{R}^n\right\}, \quad (9)$$

where t is any real number and

$$\mathcal{D}^{1,2}(\mathbb{R}^n) = \{ v \in L^{2^*}(\mathbb{R}^n) : Dv \in L^2(\mathbb{R}^n) \}, \quad \frac{1}{2^*} = \frac{1}{2} - \frac{1}{n}.$$

**Theorem 2.** Let  $\Pi = \{x \in \mathbb{R}^n : x \cdot \eta = 0\}$ . Then the following holds for a.e.  $\eta \in S^{n-1}$ : There is a rich family  $\mathcal{R}$  of Borel subsets of  $\Omega$  such that for every  $\mathcal{B} \in \mathcal{R}$  satisfying  $\mathcal{B} \subset \subset \Omega$ , the sequence of functionals

$$J_{\varepsilon}(u, \mathcal{B}) = \int_{\Omega} |Du|^2 dx + F_{\psi_{\varepsilon}}(u, \mathcal{B})$$

 $\Gamma$ -converges in the weak topology of  $H^1_0(\Omega)$  to

$$J_0(u,\mathcal{B}) = \int_{\Omega} |Du|^2 dx + \int_{\Pi \cap \mathcal{B}} \left( \int g^{\lambda}(u(x)) d\lambda \right) d\sigma(x).$$
(10)

In particular, the sequence of minimizers  $u_{\varepsilon}$  of  $J_{\varepsilon}$  converges weakly in  $H_0^1(\Omega)$  to the minimizer u of  $J_0$ .

On the right hand side of (10),  $\sigma$  denotes surface measure on  $\Pi$ .

#### 1.4 Related Results

In the paper [6] a problem similar to the present one was solved. In [6] the obstacle is given by

 $\psi \chi_{\Pi_{\varepsilon}},$ 

where  $\psi$  is a fixed function and  $\Pi_\varepsilon$  is the intersection between a hyper-plane  $\Pi$  and the set

$$\bigcup_{k\in\mathbb{Z}^n} \{a_{\varepsilon}T + \varepsilon k\},\$$

where T is a fixed subset of the unit ball. Thus in both problems the obstacle is defined on the intersection between the hyper-plane  $\Pi$  and a neighborhood of size  $a_{\varepsilon}$  of the lattice points  $\{\varepsilon k\}_{k\in\mathbb{Z}^n}$ . It is a crucial part of the problem to estimate the number of lattice points at a given distance from a subset of  $\Pi$ . For the necessary results in this direction, which come from the theory of uniform distribution, we refer to [6].

However, a main difference between the present problem and that of [6] is that the obstacle in (2) varies on a much smaller scale, of size  $a_{\varepsilon}$ . For this reason the techniques used in [6] (essentially those developed in [2]) are not fit to deal with this problem. Instead we use the methods of [3], which are more adapted to the situation at hand.

## 2 Proofs

We start by establishing some continuity properties of a certain approximation of the function  $g^{\lambda}$  in (9), that appears naturally in the proof of Theorem 2.

Lemma 1. Let

$$g_R^{\lambda}(t) = \min\left\{ \int_{B_R} |Dv|^2 dx : v - t \in H_0^1(B_R), \ v \ge \psi^{\lambda} \ q.e. \ on \ B_R \right\}.$$
(11)

Assume  $|\psi| \leq A$  and that  $\psi$  has modulus of continuity  $\rho$   $(|\psi(x) - \psi(y)| \leq \rho(|x - y|))$ . Then  $\lim_{R\to\infty} g_R^{\lambda}(t) = g^{\lambda}(t)$  and for any  $2 \leq R_0 < R_1 \leq \infty$  and any  $\lambda \in \mathbb{R}$ ,

$$|g_{R_1}^{\lambda}(t) - g_{R_2}^{\lambda}(t)| \le C(A - t)_+^2 (R_0^{2-n} - R_1^{2-n}),$$
(12)

and

$$|g_R^{\lambda+\delta}(t) - g_R^{\lambda}(t)| \le C_1 (A-t)_+^2 ((R-\delta)^{2-n} - R^{2-n}) + C_2 \rho(\delta),$$
(13)

where  $C, C_1, C_2$  depend only on n.

*Proof.* We may assume  $t \leq A$ , for otherwise  $g_R^{\lambda}(t) = 0$ . Let  $K^{\lambda}$  and  $K_R^{\lambda}$  be the set of constraints appearing in the definition of  $g^{\lambda}$  and  $g_R^{\lambda}$  respectively. That is,

$$K^{\lambda} = \left\{ v - t \in \mathcal{D}^{1,2}(\mathbb{R}^n), \ v \ge \psi^{\lambda} \text{ q.e. on } \mathbb{R}^n \right\}$$

and

$$K_R^{\lambda} = \left\{ v - t \in H_0^1(B_R), \ v \ge \psi^{\lambda} \text{ q.e. on } B_R \right\}.$$

Since  $K_{R_0}^{\lambda} \subset K_{R_1}^{\lambda} \subset K^{\lambda}$  for  $R_0 < R_1$ , we immediately obtain  $g^{\lambda}(t) \leq g_{R_1}^{\lambda}(t) \leq g_{R_0}^{\lambda}(t)$ . The claim  $\lim_{R\to\infty} g_R^{\lambda}(t) = g^{\lambda}(t)$  follows from the fact that  $C_c^{\infty}(\mathbb{R}^n)$  is dense in  $D^{1,2}(\mathbb{R}^n)$ .

Fix a smooth cut-off function  $\zeta$  with compact support in  $B_2$  such that  $\zeta \equiv 1$  on  $B_1$ . Then  $(A - t)\zeta + t \in K_R^{\lambda}$  for any  $R \geq 2, \lambda \in \mathbb{R}$  and any  $t \leq A$ . Thus

$$g_R^{\lambda}(t) \le (A-t)^2 \int_{B_2} |D\zeta|^2 dx \le C(A-t)_+^2.$$
 (14)

Let  $v \in K^{\lambda}$  satisfy  $\int_{\mathbb{R}^n} |Dv|^2 dx = g^{\lambda}(t)$ , and let  $v_R \in K_R^{\lambda}$  satisfy  $\int_{B_R} |Dv_R|^2 dx = g_R^{\lambda}(t)$ . To estimate  $v - v_R$  we construct a barrier h that is the solution to  $\Delta h = 0$  in  $\mathbb{R}^n \setminus B_1$ ,  $h - t \in \mathcal{D}^{1,2}(\mathbb{R}^n)$  and h = A on  $B_1$ . In  $\mathbb{R}^n \setminus B_1$ , h - v is harmonic, on  $B_1$ ,  $h - v \ge 0$  and  $h - v \to 0$  at infinity. It follows from the maximum principle that  $v \le h$  in  $\mathbb{R}^n$ . The function h is spherically symmetric and has the explicit expression

$$h(r) = (A - t)r^{2-n} + t,$$

for r > 1, where r = |x|. It follows that

$$v(x) \le (A-t)R^{2-n} + t$$
 on  $\mathbb{R}^n \setminus B_R$ .

Thus

$$\hat{v}_R = \max(t, v - (1 - \zeta)(A - t)R^{2-n})$$

belongs to  $K_R^{\lambda}$ . Hence

$$\begin{split} g_{R}^{\lambda}(t) &\leq \int_{B_{R}} |D\hat{v}_{R}|^{2} dx \\ &\leq \int_{B_{R}} |Dv|^{2} dx + 2(A-t)R^{2-n} \int_{B_{R}} D\zeta Dv dx + ((A-t)R^{2-n})^{2} \int_{B_{R}} |D\zeta|^{2} dx \\ &\leq g^{\lambda}(t) + 2(A-t)R^{2-n} \|D\zeta\|_{L^{2}(B_{R})} \sqrt{g^{\lambda}(t)} + ((A-t)R^{2-n})^{2} \int_{B_{R}} |D\zeta|^{2} dx. \end{split}$$

Hence we obtain, using (14),

$$|g^{\lambda}(t) - g^{\lambda}_{R}(t)| \le C(A - t)^{2} R^{2-n}.$$
(15)

If  $2 < R_0 < R_1$ , we find in a similar way that

$$v_{R_1} \le h_{R_1} = (A-t)\frac{r^{2-n} - R_1^{2-n}}{1 - R_1^{2-n}} + t \text{ on } B_{R_1} \setminus B_1,$$

and that

$$\hat{v}_{R_0} = \max\left(t, v_{R_1} - (1-\zeta)(A-t)\frac{R_0^{2-n} - R_1^{2-n}}{1 - R_1^{2-n}}\right)$$

belongs to  $K_{R_0}^{\lambda}$ . From this we obtain the estimate

$$|g_{R_1}^{\lambda}(t) - g_{R_2}^{\lambda}(t)| \le C(A-t)^2 (R_0^{2-n} - R_1^{2-n}).$$
(16)

Next we prove the continuity w.r.t.  $\lambda$ . For any  $\gamma > 0$  there exists a  $\delta > 0$   $(\delta = \rho^{-1}(\gamma))$  such that

$$\psi^{\lambda}(x+\delta\eta) - \gamma < \psi^{\lambda+\delta}(x) \le \psi^{\lambda}(x+\delta\eta) + \gamma.$$

Let

$$h_R = \frac{r^{2-n} - R^{2-n}}{1 - R^{2-n}},$$

for r = |x| > 1,  $h_R = 1$  on  $B_1$ . Let  $v_{R-\delta}^{\lambda} \in K_{R-\delta}^{\lambda}$  satisfy  $\int_{B_{R-\delta}} |Dv_{R-\delta}^{\lambda}|^2 dx = g_{R-\delta}^{\lambda}$ . Then  $w_R(x) = v_{R-\delta}^{\lambda}(x + \delta\eta) + \gamma h_R(x)$  belongs to  $K_R^{\lambda+\delta}$ . Hence,

$$\begin{split} g_{R}^{\lambda+\delta}(t) &\leq \int_{B_{R}} |Dw_{R}|^{2} dx \\ &= \int_{B_{R}} |Dv_{R-\delta}^{\lambda}(x+\delta\eta)|^{2} dx + \gamma^{2} \int_{B_{R}} |Dh_{R}|^{2} dx + 2\gamma \int_{B_{R}} Dh_{R} Dv_{R-\delta}^{\lambda} dx \\ &\leq g_{R}^{\lambda}(t) + C(A-t)^{2} ((R-\delta)^{2-n} - R^{2-n}) \\ &+ \gamma^{2} \int_{B_{R}} |Dh_{R}|^{2} dx + 2\gamma \|Dv_{R-\delta}^{\lambda}\|_{L^{2}(B_{R})} \|Dh_{R}\|_{L^{2}(B_{R})}. \end{split}$$

It is easy to check that  $\int_{B_R} |Dh_R|^2 dx$  is bounded uniformly in R. In fact, as  $R \to \infty$ ,  $\int_{B_R} |Dh_R|^2 dx \to \operatorname{cap}(B_1)$ , the capacity of the unit ball. By interchanging the roles of  $g_R^{\lambda+\delta}(t)$  and  $g_R^{\lambda}(t)$  we obtain a lower bound on  $g_R^{\lambda+\delta}(t) - g_R^{\lambda}(t)$ . Thus for any  $\gamma > 0$ , we have (assuming  $\gamma < 1$ )

$$|g_R^{\lambda+\delta}(t) - g_R^{\lambda}(t)| \le C_1 (A-t)^2 ((R-\delta)^{2-n} - R^{2-n}) + C_2 \gamma.$$
(17)

We now turn to the

proof of Theorem 2. Let  $w_{\varepsilon}^k$  be the solution to

$$\min\left\{\int_{Q_{\varepsilon}^{k}} |Dw|^{2} dx : w \ge \psi_{\varepsilon} \text{ q.e. on } Q_{\varepsilon}^{k}, w = t \text{ on } Q_{\varepsilon}^{k} \setminus B_{\varepsilon/2}^{k}\right\}.$$
 (18)

The following definition will be important in the sequel. In order to simplify notation we set  $P = \Pi \cap \mathcal{B}$ .

**Definition 2.** Let  $\lambda_{\varepsilon}^k$  be the unique real number such that

 $\begin{aligned} Q_{\varepsilon}^k \cap P &= Q_{\varepsilon} \cap \{P + \lambda_{\varepsilon}^k \eta\} \pmod{\varepsilon}, \quad \text{if } Q_{\varepsilon}^k \cap P \neq \emptyset. \end{aligned}$  If  $Q_{\varepsilon}^k \cap P &= \emptyset \text{ we set } \lambda_{\varepsilon}^k = \infty. \end{aligned}$ 

Let  $y = x - \varepsilon k$ . Then

$$y + \varepsilon k \in Q_{\varepsilon}^k \cap P \Longleftrightarrow y \in Q_{\varepsilon} \cap \{P + \lambda_{\varepsilon}^k \eta\}.$$

Thus

$$\begin{split} &\int_{Q_{\varepsilon}^{k}} |Dw_{\varepsilon}^{k}|^{2} dx \\ &= \min \left\{ \int_{Q_{\varepsilon}} |Dw|^{2} dx : w \geq \psi_{\varepsilon}^{\lambda_{\varepsilon}^{k}} \text{ q.e. on } Q_{\varepsilon}, \ w = t \text{ on } Q_{\varepsilon} \setminus B_{\varepsilon/2} \right\}, \end{split}$$

where  $\psi_{\varepsilon}^{\lambda_{\varepsilon}^{k}}$  is  $\psi_{\varepsilon}$  with  $P + \lambda_{\varepsilon}^{k}\eta$  in place of P. Clearly,  $w_{\varepsilon}^{k} = t$  if  $\psi_{\varepsilon}^{\lambda_{\varepsilon}^{k}} \leq t$ . In particular,  $w_{\varepsilon}^{k} = t$  if  $Q_{\varepsilon}^{k} \cap (\Omega \cap P) = \emptyset$ . Let  $z = a_{\varepsilon}^{-1}y$ . Then, noting that  $a_{\varepsilon}z = y \in Q_{\varepsilon} \cap \{P + \lambda_{\varepsilon}^{k}\eta\} \iff z \in Q_{\varepsilon/a_{\varepsilon}} \cap \{P + (\lambda_{\varepsilon}^{k}/a_{\varepsilon})\eta\},$ 

$$\int_{Q_{\varepsilon}^{k}} |Dw_{\varepsilon}^{k}|^{2} dx = \min \left\{ a_{\varepsilon}^{n-2} \int_{Q_{\varepsilon/a_{\varepsilon}}} |Dw|^{2} dx : w \ge \psi^{\lambda_{\varepsilon}^{k}/a_{\varepsilon}} \text{ q.e. on } Q_{\varepsilon/a_{\varepsilon}}, \right.$$
  
and  $w = t$  on  $Q_{\varepsilon/a_{\varepsilon}} \setminus B_{\varepsilon/2a_{\varepsilon}} \left. \right\}.$ 

Let  $R_{\varepsilon} = \varepsilon/2a_{\varepsilon}$ . The choice of  $a_{\varepsilon}$  implies that  $\lim_{\varepsilon \to 0} R_{\varepsilon} = \infty$ . Since w - t has its support in  $B_{R_{\varepsilon}}$  and  $\psi^{\lambda_{\varepsilon}^k/a_{\varepsilon}} = -\infty$  outside  $B_1 \subset B_{R_{\varepsilon}}$ , we have

$$\min\left\{a_{\varepsilon}^{n-2}\int_{Q_{\varepsilon/a_{\varepsilon}}}|Dw|^{2}\,dx:w\geq\psi^{\lambda_{\varepsilon}^{k}/a_{\varepsilon}}\text{ q.e. on }Q_{\varepsilon/a_{\varepsilon}},\right.$$
  
and  $w=t$  on  $Q_{\varepsilon/a_{\varepsilon}}\setminus B_{\varepsilon/2a_{\varepsilon}}\right\}=$ 

$$= \min \left\{ a_{\varepsilon}^{n-2} \int_{B_{R_{\varepsilon}}} |Dw|^2 \, dx : w \ge \psi^{\lambda_{\varepsilon}^k/a_{\varepsilon}} \text{ q.e. on } B_{R_{\varepsilon}}, \right.$$
  
and  $w - t \in H_0^1(B_{R_{\varepsilon}}) \right\}$ 
$$= a_{\varepsilon}^{n-2} g_{R_{\varepsilon}}^{\lambda_{\varepsilon}^k/a_{\varepsilon}}(t).$$

It is clear that  $\psi^{\lambda_{\varepsilon}^k/a_{\varepsilon}} \equiv -\infty$  for small enough  $\varepsilon > 0$  if  $a_{\varepsilon} = o(\lambda_{\varepsilon})$ . Choose  $\lambda_0 < \lambda_1$  such that  $B_1 \cap \{P + \lambda\eta\} = \emptyset$  if  $\lambda \notin [\lambda_0, \lambda_1]$ . Let  $\delta > 0$  be a small number such that  $\lambda_1 = \lambda_0 + M\delta$  for some positive integer M, and let  $\lambda_j = \lambda_0 + j\delta$ . Now set  $\lambda_{\varepsilon,j} = a_{\varepsilon}\lambda_j$  and let

$$I_{\varepsilon,j} = \{Q_{\varepsilon} \cap \{P + \lambda\eta\} : \lambda_{\varepsilon,j} \le \lambda \le \lambda_{\varepsilon,j+1}\},\$$
  
$$I_{\varepsilon,j}^{k} = \{I_{\varepsilon,j} + \varepsilon k\}, \quad k \in \mathbb{Z}^{n}.$$

Let  $A_{\varepsilon,j}$  be the number of  $k \in \mathbb{Z}^n$  for which P and  $I_{\varepsilon,j}^k$  has non-empty intersection. This is precisely the number of  $k = (k', k_n)$  such that  $\varepsilon k_n$  and  $\alpha \varepsilon k'$  belong to the same cube  $Q_{\varepsilon}^k$ , and  $\lambda_{\varepsilon}^k \in I_{\varepsilon,j}$ , where we use the notation in (1). Let

$$P_{\varepsilon} = \{ k \in \mathbb{Z}^n : Q_{\varepsilon}^k \cap P \neq \emptyset \}.$$

Thus if

$$\mathbb{K}_{\varepsilon,j} = \{ k \in P_{\varepsilon} : \lambda_{\varepsilon}^k \in I_{\varepsilon,j} \},\$$

then

$$A_{\varepsilon,j} = \# \mathbb{K}_{\varepsilon,j}$$

It was proven in [6], Lemma 5.2.2, that for a.e.  $\eta \in S^{n-1}$ ,

$$A_{\varepsilon,j} = |P| \frac{\delta a_{\varepsilon}}{\varepsilon^n} + o(a_{\varepsilon} \varepsilon^{-n}).$$
(19)

To make the statement more precise we introduce

$$N_{\varepsilon} = \#\{k' \in \mathbb{Z}^{n-1} \cap \operatorname{proj}_{\mathbb{R}^{n-1}} \varepsilon^{-1} P\}.$$

Then, since the intersection between P and  $I_{\varepsilon,j}^k$  is completely determined by the value of  $\varepsilon \alpha k'$  at a point  $(\varepsilon k', \alpha \varepsilon k') \in P$ , we have

$$A_{\varepsilon,j} = \# \left\{ k' \in \mathbb{Z}^{n-1} \cap \operatorname{proj}_{\mathbb{R}^{n-1}} \varepsilon^{-1} P : \alpha k' / \mathbb{Z} \in [p_j, p_j + \delta a_{\varepsilon} / (\eta_n \varepsilon)] / \mathbb{Z} \right\},\$$

where  $p_j$  is chosen such that

$$P \cap I_{\varepsilon,j}^k \neq \emptyset$$
 iff  $\alpha k'/\mathbb{Z} \in [p_j, p_j + \delta a_{\varepsilon}/(\eta_n \varepsilon)]/\mathbb{Z}$ .

Note that the distance  $\delta a_{\varepsilon}$  in  $\eta$  (normal) direction between two planes, corresponds to the distance  $\delta a_{\varepsilon}/\eta_n$  in  $e_n$  direction between these planes. Using tools from the theory of uniform distribution mod 1, it can be shown that

$$\left|\frac{A_{\varepsilon,j}}{N_{\varepsilon}} - \frac{\delta a_{\varepsilon}}{\varepsilon \eta_n}\right| = o(\varepsilon^s), \quad \text{for any } s \in (0,1).$$

This implies (19) since  $a_{\varepsilon}/\varepsilon \geq \sqrt{\varepsilon}$  for  $n \geq 3$ . Define  $w_{\varepsilon}$  by  $w_{\varepsilon} = w_{\varepsilon}^{k}$  on  $Q_{\varepsilon}^{k}$ . Since  $w_{\varepsilon}^{k} = t$  on  $\partial B_{r_{\varepsilon}}^{k}$ ,  $w_{\varepsilon} \in H^{1}(\Omega)$  and, noting that  $w_{\varepsilon}^{k} \equiv t$  if  $k \notin \mathbb{K}_{\varepsilon,j}$  for some j,

$$\int_{\Omega} |Dw_{\varepsilon}|^2 dx = \sum_{j=0}^{M} \sum_{k \in \mathbb{K}_{\varepsilon,j}} \int |Dw_{\varepsilon}^k|^2 dx$$
(20)

$$= \sum_{j=0}^{M} \sum_{k \in \mathbb{K}_{\varepsilon,j}} a_{\varepsilon}^{n-2} \left( g_{R_{\varepsilon}}^{\lambda_{\varepsilon}^{k}/a_{\varepsilon}}(t) - g_{R_{\varepsilon}}^{\lambda_{j}}(t) \right) + \sum_{j=0}^{M} a_{\varepsilon}^{n-2} A_{\varepsilon,j} g_{R_{\varepsilon}}^{\lambda_{j}}(t).$$
(21)

Since  $|\lambda_{\varepsilon}^k/a_{\varepsilon} - \lambda_j| \leq \delta$  when  $k \in \mathbb{K}_{\varepsilon,j}$ , we have for such k that

$$\left|g_{R_{\varepsilon}}^{\lambda_{\varepsilon}^{k}/a_{\varepsilon}}(t) - g_{R_{\varepsilon}}^{\lambda_{j}}(t)\right| \leq C_{1}(A-t)^{2}_{+}((R_{\varepsilon}-\delta)^{2-n} - R_{\varepsilon}^{2-n}) + C_{2}\rho(\delta) =: E(\varepsilon,\delta),$$

by (13) in Lemma 1. Hence the first term in (21) is bounded by

$$\sum_{j=0}^{M} A_{\varepsilon,j} a_{\varepsilon}^{n-2} E(\varepsilon,\delta) \le C \sum_{j=0}^{M} |P| \delta \frac{a_{\varepsilon}^{n-1}}{\varepsilon^{n}} E(\varepsilon,\delta) \le C |P| E(\varepsilon,\delta),$$
(22)

where we used (19), the fact that  $a_{\varepsilon}^{n-1}/\varepsilon^n = 1$  by the choice of  $a_{\varepsilon}$  in (3) and that  $M = 1/\delta$ . The right hand side of (22) clearly tends to zero as  $\varepsilon, \delta \to 0$ in any order. The term  $a_{\varepsilon}^{n-2}A_{\varepsilon,j} g_{R_{\varepsilon}}^{\lambda_j}(t)$  converges to  $|P|\delta g^{\lambda_j}(t)$  as  $\varepsilon \to 0$ . Hence,

$$\int |Dw_{\varepsilon}|^2 dx = \sum_{j=0}^M \sum_{k \in \mathbb{K}_{\varepsilon,j}} \int |Dw_{\varepsilon}^k|^2 dx = O(\rho(\delta)) + \sum_{j=0}^M A_{\varepsilon,j} g_{R_{\varepsilon}}^{\lambda_j}(t)$$
$$\to \sum_{j=0}^M \delta |P| g^{\lambda_j}(t),$$

as  $\varepsilon \to 0$ . Letting  $\delta \to 0$ , we obtain

$$\int_{\Omega} |Dw_{\varepsilon}|^2 dx = \sum_k \int_{\Omega} |Dw_{\varepsilon}^k|^2 dx \to |P| \int_{\lambda_0}^{\lambda_1} g^{\lambda}(t) d\lambda.$$
(23)

The next step is to show that  $w_{\varepsilon} \rightharpoonup t$  in  $H^1(\Omega)$ . Since  $w_{\varepsilon} - t \in H_0(B^k_{\varepsilon/2})$ , Poincare's inequality implies that

$$\int_{B_{\varepsilon/2}^k} |w_{\varepsilon}^k - t|^2 dx \le \varepsilon \int_{B_{\varepsilon/2}^k} |Dw_{\varepsilon}^k|^2 dx.$$

Indeed, the Poincare constant of a ball of radius R does not exceed R. Thus

$$\int_{\Omega} |w_{\varepsilon} - t|^2 dx = \sum_{k} \int_{B_{\varepsilon/2}^k} |w_{\varepsilon}^k - t|^2 dx$$
(24)

$$\leq \varepsilon \sum_{k} \int_{B_{\varepsilon/2}^{k}} |Dw_{\varepsilon}^{k}|^{2} dx = \varepsilon^{2} \int_{\Omega} |Dw_{\varepsilon}|^{2} dx.$$
(25)

By (23)  $\{w_{\varepsilon}\}_{\varepsilon}$  is bounded in  $H_0^1(\Omega)$  and hence has a weakly convergent subsequence quence. From (24)-(25) it follows that every weakly convergent subsequence must converge to t, thus the entire sequence  $\{w_{\varepsilon}\}_{\varepsilon}$  converges weakly to t.

By Theorem 1,  $J_{\varepsilon}(u) = \int_{\Omega} |Du|^2 dx + F_{\psi_{\varepsilon}}(u, \mathcal{B})$  has a subsequence that  $\Gamma$ -converges to a functional of the type  $J_0(u) = \int_{\Omega} |Du|^2 dx + \int_{\mathcal{B}} f(x, u) d\mu$ . We will prove that for each  $t \in \mathbb{R}$ ,

$$\int_{\mathcal{B}} f(x,t)d\mu = \sigma(\Pi \cap \mathcal{B}) \int g^{\lambda}(t)d\lambda.$$
(26)

Let us show that the theorem follows from (26). Due to (26) and the fact that the family of sets  $\mathcal{R} \ni \mathcal{B}$  is dense in the Borel subsets of  $\Omega$ ,  $f(x,t)d\mu$  is a measure on  $\Pi$ , absolutely continuous w.r.t.  $\sigma$ . Hence  $f(x,t)d\mu = h(x,t)d\sigma$ for some  $h(x,t) \in L^1_{loc}(\Pi,\sigma)$ . But

$$\int_{\Pi \cap \mathcal{B}} h(x,t) d\sigma = \sigma(\Pi \cap \mathcal{B}) \int g^{\lambda}(t) d\lambda$$

for all  $t \in \mathbb{R}$  and all  $\mathcal{B} \in \mathcal{R}$  implies that h is independent of x, thus  $h(x,t) = h(t) = \int g^{\lambda}(t) d\lambda$ .

We now prove (26). Choose  $v \in C_c^{\infty}(\Omega)$  such that v = t on a neighbourhood of  $\mathcal{B}$ . Let

$$v_{\varepsilon}(x) = \begin{cases} w_{\varepsilon}(x), & \text{if } x \in \mathcal{B}, \\ v(x), & \text{if } x \in \Omega \setminus \mathcal{B}. \end{cases}$$
(27)

Then clearly  $v_{\varepsilon} \rightharpoonup v$  in  $H^1(\Omega)$ . According to Definition 1 (i),

$$\int_{\Omega} |Dv|^2 dx + \int_{\mathcal{B}} f(u, x) d\mu = \int_{\Omega \setminus \mathcal{B}} |Dv|^2 dx + \int_{\mathcal{B}} f(t, x) d\mu$$
  
$$\leq \liminf_{\varepsilon \to 0} \int_{\Omega} |Dv_{\varepsilon}|^2 dx = \int_{\Omega \setminus \mathcal{B}} |Dv|^2 dx + \sigma(\mathcal{B} \cap \Pi) \int g^{\lambda}(t) d\lambda.$$

It remains to prove that

$$\int_{\mathcal{B}} f(x,t)d\mu \ge \sigma(\mathcal{B} \cap \Pi)g^{\lambda}(t)d\lambda.$$
(28)

Let  $z_{\varepsilon}$  be a sequence given by Definition 1 (ii), i.e.  $z_{\varepsilon} \rightharpoonup v$  and  $\limsup_{\varepsilon} J_{\varepsilon}(z_{\varepsilon}) \leq J_0(v)$ . By (i) in the same definition, we have  $\lim_{\varepsilon \to 0} J_{\varepsilon}(z_{\varepsilon}) = J_0(v)$ . Since v is bounded we may assume  $z_{\varepsilon}$  is bounded. To see this we assume  $|v| \leq C$  and claim that

$$\bar{z}_{\varepsilon} = \min(z_{\varepsilon}^+, 2C) - \min(z_{\varepsilon}^-, 2C) \rightharpoonup v.$$

Indeed,  $\bar{z}_{\varepsilon}$  is uniformly bounded in  $H^1(\Omega)$  and therefore has a weak limit in this space. Moreover,

$$\int_{\Omega} |\bar{z}_{\varepsilon} - v|^2 dx = \int_{\Omega \setminus \{|z_{\varepsilon}| > 2C\}} |z_{\varepsilon} - v|^2 dx - \int_{\{z_{\varepsilon} > 2C\}} |2C - v|^2 dx$$
$$- \int_{\{z_{\varepsilon} < -2C\}} |-2C - v|^2 dx.$$

Since  $z_{\varepsilon} \to v$  strongly in  $L^2(\Omega)$  and

$$\int_{\Omega} |z_{\varepsilon} - v|^2 dx \ge C^2 \text{measure}(\{|z_{\varepsilon}| > 2C\}),$$

measure({ $|z_{\varepsilon}| > 2C$ })  $\rightarrow 0$  and hence  $\bar{z}_{\varepsilon} \rightarrow v$  strongly in  $L^{2}(\Omega)$ . Additionally,  $\int |D\bar{z}_{\varepsilon}|^{2} dx \leq \int |Dz_{\varepsilon}|^{2} dx$ , which implies, again by (i) in Definition 1,

$$\lim_{\varepsilon \to 0} J_{\varepsilon}(\bar{z}_{\varepsilon}) = J_0(v) = \int_{\Omega \setminus \mathcal{B}} |Dv|^2 dx + \int_{\mathcal{B}} f(t, x) d\mu.$$

Thus if we let  $v_{\varepsilon}$  be the function given by (27), (28) follows if we prove

$$\begin{cases} \lim_{\varepsilon \to 0} \int_{\Omega} |Dv_{\varepsilon}|^2 dx \leq \lim_{\varepsilon \to 0} \int_{\Omega} |Dz_{\varepsilon}|^2 dx, \\ \text{for all } z_{\varepsilon} \in H_0^1(\Omega) \text{ such that } z_{\varepsilon} \geq \psi_{\varepsilon}, \\ z_{\varepsilon} \rightharpoonup v \text{ and } \sup_{\varepsilon > 0} \|z_{\varepsilon}\|_{L^{\infty}} < \infty. \end{cases}$$
(29)

By convexity of the functional  $v \mapsto \int_{\Omega} |Dv|^2 dx$ , we have

$$\int_{\Omega} |Dz_{\varepsilon}|^2 - |Dv_{\varepsilon}|^2 dx \ge 2 \int_{\Omega} Dv_{\varepsilon} (Dz_{\varepsilon} - Dv_{\varepsilon}) dx \tag{30}$$

$$= \langle -\Delta v_{\varepsilon}, z_{\varepsilon} - v_{\varepsilon} \rangle = \int_{\Omega \setminus \mathcal{B}} -\Delta v(z_{\varepsilon} - v)dx + \sum_{k} \langle -\Delta w_{\varepsilon}^{k}, z_{\varepsilon} - w_{\varepsilon}^{k} \rangle, \quad (31)$$

where the sum is taken over

$$\{k \in \mathbb{Z}^n : \Pi \cap \mathcal{B} \subset \{a_{\varepsilon}\{y : \psi(y) > -\infty\} + \varepsilon k\} \ (\subset B^k_{a_{\varepsilon}/2})\}.$$

The first term in (31) goes to zero since v is smooth and  $z_{\varepsilon} \rightharpoonup v$ . The Laplacian of  $w_{\varepsilon}^k$  consists of two measures  $\mu_{\varepsilon}^k$  and  $\nu_{\varepsilon}^k$  such that

$$-\Delta w_{\varepsilon} = \mu_{\varepsilon}^k - \nu_{\varepsilon}^k,$$

where

$$\nu_{\varepsilon}^{k}(E) = -\int_{E \cap Q_{\varepsilon}^{k}} \frac{\partial w_{\varepsilon}^{k}}{\partial n} dS,$$

and

$$\operatorname{supp}\mu_{\varepsilon}^{k} \subset \{w_{\varepsilon}^{k} = \psi^{\varepsilon}\} \subset B_{a_{\varepsilon}}^{k}, \tag{32}$$

which follows from the fact that  $w_{\varepsilon}^k$  solves (18) (see [5]). From (32) and the fact that  $z_{\varepsilon} \geq \psi_{\varepsilon}$  it follows that

$$\int_{Q_{\varepsilon}^{k}} (z_{\varepsilon} - w_{\varepsilon}^{k}) d\mu_{\varepsilon}^{k} = \int_{Q_{\varepsilon}^{k}} (z_{\varepsilon} - \psi_{\varepsilon}) d\mu_{\varepsilon}^{k} + \int_{Q_{\varepsilon}^{k}} (\psi_{\varepsilon} - w_{\varepsilon}^{k}) d\mu_{\varepsilon}^{k}$$
$$= \int_{Q_{\varepsilon}^{k}} (z_{\varepsilon} - \psi_{\varepsilon}) d\mu_{\varepsilon}^{k} \ge 0.$$

It remains to show that

$$\lim_{\varepsilon \to 0} \sum_{k} \int_{Q_{\varepsilon}^{k}} (z_{\varepsilon} - w_{\varepsilon}^{k}) d\nu_{\varepsilon}^{k} = 0.$$

Let  $W^k_\varepsilon$  solve

$$\min\left\{\int_{Q_{\varepsilon}^{k}}|DW|^{2}dx:W-t\in H_{0}^{1}(B_{\varepsilon/2}^{k})\text{ and }W\geq\max\psi=A\text{ on }B_{a_{\varepsilon}}^{k}\right\}.$$

Since  $W_{\varepsilon}^{k} = w_{\varepsilon}^{k}$  on  $\partial B_{\varepsilon/2}^{k}$ ,  $W_{\varepsilon}^{k} \ge w_{\varepsilon}^{k}$  on  $B_{a_{\varepsilon}}^{k}$  and  $W_{\varepsilon}^{k}$  are harmonic in  $B_{\varepsilon/2}^{k} \setminus B_{a_{\varepsilon}}^{k}$ , we get  $W_{\varepsilon}^{k} \ge w_{\varepsilon}^{k}$  in  $B_{\varepsilon/2}^{k}$  from the maximum principle, hence

$$-\frac{\partial W_{\varepsilon}^k}{\partial n} \geq -\frac{\partial w_{\varepsilon}^k}{\partial n} \text{ on } \partial B_{\varepsilon/2}^k.$$

Thus if we let

$$\hat{\nu}_{\varepsilon}^{k}(E) = \int_{\partial B_{\varepsilon/2}^{k} \cap E} -\frac{\partial W_{\varepsilon}^{k}}{\partial n} dS,$$

and set  $\hat{\nu}_{\varepsilon} = \sum_{k} \hat{\nu}_{\varepsilon}^{k}$ ,  $\nu_{\varepsilon} = \sum_{k} \nu_{\varepsilon}^{k}$ , then  $\hat{\nu}_{\varepsilon} \ge \nu_{\varepsilon}$ . In [6] (see the proof of Lemma 2.0.8 therein) it was shown that

$$\lim_{\varepsilon \to 0} \int_{\Omega} (h_{\varepsilon} - h) d\hat{\nu}_{\varepsilon} = 0, \qquad (33)$$

whenever  $h_{\varepsilon} \rightharpoonup h$  in  $H_0^1(\Omega)$  and  $\sup_{\varepsilon>0} \|h_{\varepsilon}\|_{L^{\infty}} < \infty$ . Since  $\nu_{\varepsilon} \leq \hat{\nu}_{\varepsilon}$ , it follows that (33) holds for  $\nu_{\varepsilon}$  after writing  $(h_{\varepsilon} - h) = (h_{\varepsilon} - h)_+ - (h_{\varepsilon} - h)_-$ . This proves (29). Since the  $\Gamma$ -limit  $J_0$  does not depend on the particular  $\Gamma$ -convergent subsequence, the entire sequence  $J_{\varepsilon}$   $\Gamma$ -converges to  $J_0$ .

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