

# PERIODIC NAVIE-STOKES SOLUTIONS WITH CRYSTALLOGRAPHIC SYMMETRIES

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We study periodic solutions to the Navier-Stokes equations (NSE) with arbitrary vector periods  $\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3$ . Let  $\Lambda$  be the lattice of periods  $\mathbf{p} = n_1\mathbf{p}_1 + n_2\mathbf{p}_2 + n_3\mathbf{p}_3$ ,  $n_i \in \mathbb{Z}$ , and  $\Lambda^*$  be the reciprocal lattice of vectors  $\mathbf{k} = n_1\mathbf{k}_1 + n_2\mathbf{k}_2 + n_3\mathbf{k}_3$  where  $\mathbf{k}_i \cdot \mathbf{p}_j = 2\pi\delta_{ij}$ . Functions  $\mathbf{V}(t, \mathbf{x}) = \sum \mathbf{V}_{\mathbf{k}}(t) \exp(i\mathbf{k} \cdot \mathbf{x})$ ,  $\nabla p(t, \mathbf{x}) = \mathbf{p}_0(t) + \sum p_{\mathbf{k}}(t) \exp(i\mathbf{k} \cdot \mathbf{x})$  are periodic with periods  $\mathbf{p}_1, \mathbf{p}_2, \mathbf{p}_3$ . Here  $\mathbf{k} \in \Lambda^*$ ,  $\mathbf{V}_{-\mathbf{k}} = \overline{\mathbf{V}_{\mathbf{k}}}$ ,  $p_{-\mathbf{k}} = \overline{p_{\mathbf{k}}}$ , and  $\mathbf{V}_{\mathbf{k}} \cdot \mathbf{k} = 0$ . For periodic solutions, the NSE reduce to the dynamical system

$$\dot{\mathbf{V}}_{\mathbf{n}} = -\mathbf{n}^2\nu\mathbf{V}_{\mathbf{n}} + \frac{i}{\mathbf{n}^2}\mathbf{n} \times \left( \mathbf{n} \times \sum_{\mathbf{k} \in \Lambda^*} (\mathbf{V}_{\mathbf{k}} \cdot \mathbf{n}) \mathbf{V}_{\mathbf{n}-\mathbf{k}} \right), \quad \dot{\mathbf{V}}_0 = -\frac{1}{\rho}\mathbf{p}_0, \quad (1)$$

where vectors  $\mathbf{k}, \mathbf{n} \in \Lambda^*$ . We show that dynamical systems (1) for different triples of periods  $\mathbf{p}_j$  generically are not equivalent to each other: the moduli space of non-equivalent systems (1) has dimension 6.

We construct exact NSE solutions with crystallographic symmetry groups  $G$  that have pure rotational point groups  $\Gamma \subset SO(3)$ . There are 52 such groups  $G$  among 219 nonisomorphic crystallographic groups in three dimensions. The point group  $\Gamma$  can be either cyclic  $C_n$ , or dihedral  $D_n$ ,  $n = 2, 3, 4, 6$ , or tetrahedral  $T$  or octahedral group  $O$ . The constructed exact solutions depend upon all four variables  $t, x_1, x_2, x_3$ .

We obtain complete classification of periodic solutions with pairwise non-interacting Fourier modes. For such solutions, the non-zero Fourier components  $\mathbf{V}_{\mathbf{k}}(t) \in \mathbb{C}^3$  correspond to vectors  $\mathbf{k}$  of the reciprocal lattice  $\Lambda^*$  that necessarily belong either to the spheres  $\mathbf{k}^2 = \alpha^2$ , or to the circumferences  $\mathbf{k} \cdot \mathbf{e} = 0$ ,  $\mathbf{k}^2 = \alpha^2$ , or to the straight lines  $\mathbf{k} = \lambda\mathbf{n}$  or to the planes  $\mathbf{k} \cdot \mathbf{e} = 0$ , where  $\mathbf{k}, \mathbf{n} \in \Lambda^*$  and  $\mathbf{e} \in \Lambda$ .

The system (1) has an infinite-dimensional Lie group of symmetries  $G = H(\Lambda) \dot{\times} \mathbf{A}$  where  $H(\Lambda)$  is the holohedry group of the lattice  $\Lambda$  (or  $\Lambda^*$ ) and  $\mathbf{A}$  is the abelian Lie group of vector-valued functions  $\mathbf{S}(t)$ . Applying the symmetry transforms  $\tilde{\mathbf{V}}_{\mathbf{k}} = \exp(i\mathbf{k} \cdot \mathbf{S}(t))Q\mathbf{V}_{Q^{-1}(\mathbf{k})}(t)$ ,  $\tilde{\mathbf{V}}_0 = Q\mathbf{V}_0(t) - \dot{\mathbf{S}}(t)$ ,  $\tilde{p}_{\mathbf{k}} = \exp(i\mathbf{k} \cdot \mathbf{S}(t))p_{Q^{-1}(\mathbf{k})}(t)$ ,  $\tilde{\mathbf{p}}_0 = Q\mathbf{p}_0(t) + \rho\ddot{\mathbf{S}}(t)$ , where  $Q \in H(\Lambda)$ , we

demonstrate the non-uniqueness of solutions to the Cauchy problem for the periodic NSE with  $\mathbf{V}_0(t) \neq 0$ .

**References:**

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