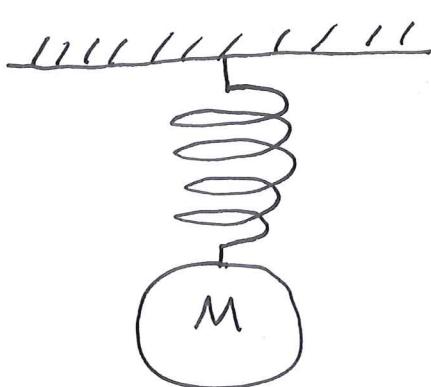


Fourier Analysis

- 1) Book Stein & Shakarovich.
- 2) Every Tuesday 1-3 pm here
- 3) Examination.
 - i) For E-D weekly assignments
(Not all will be assessed)
 - ii) FOR C-B You need a D on the assignments and write an exam
 - iii) For A you need a B and pass an oral exam.
- 4) We will follow chapter 1-6 in Stein - Shakarovich's book for the first 10 (or so) lectures.
Then we will do something else (maybe chapter 7-8).
What I really want to do is mathematics

Let M be a weight of mass m

hanging in a spring (without any gravity) such



that the position at time t is $y(t)$.

We assume that

the spring exerts a force
 $-ky(t)$ on M and

that Newton's law $my''(t) = F = -ky(t)$
applies. That is $y(t)$ satisfies the differential eq.

$$y'' + \frac{k}{m} y = 0 \quad (1).$$

where $c = \sqrt{\frac{k}{m}}$.

We know that the solution has the form

$$y(t) = a \cos(ct) + b \sin(ct), \quad \text{for some } a, b \in \mathbb{R}.$$

This is extremely simple (first year analysis).

Given some initial data

$$y(t_0) = A \quad \& \quad y'(t_0) = B$$

we can solve $y(t)$.

Example 2. (Vibrating string)

Question? ~~But~~ What equation determines the evolution of a vibrating string?



Two things seems to be of importance

1) The mass of the string. M

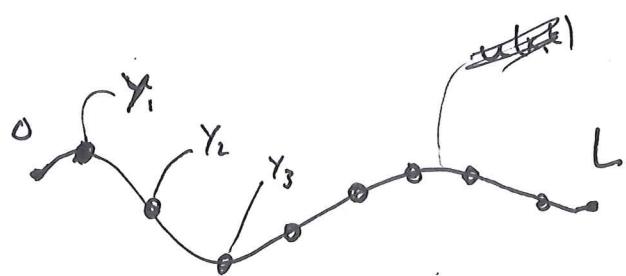
2) The tension in the string. T

Assume T to be independent of x and t .

Discretisation: Let us look at the "approximating" problem with a weightless string with N (=large number)

of masses $\frac{M}{N}$ equally distributed along the string. That is

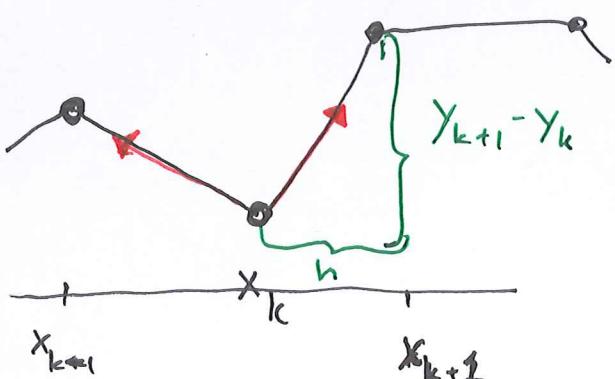
There is a mass $\frac{M}{N}$ at each point $\frac{kL}{N}$ along the string.



$$x_0 \quad x_1 = \frac{L}{N} \quad x_k = \frac{kL}{N}$$

$$\frac{L}{N} = h$$

We will call



The force acting on the y direction

(x_k, y_k) from the right is

$$T \sin(\arctan(\frac{y_{k+1} - y_k}{h})) =$$

$$= T \frac{\frac{y_{k+1} - y_k}{h}}{\sqrt{1 + \left(\frac{y_{k+1} - y_k}{h}\right)^2}} \approx \left\{ \text{Taylor} \right\}$$

$$\approx T \frac{y_{k+1} - y_k}{h} \left(1 + \frac{1}{2} \left(\frac{y_{k+1} - y_k}{h} \right)^2 \right) \approx \left\{ \begin{array}{l} \text{Assume } h \text{ is} \\ \text{that } \frac{y_{k+1} - y_k}{h} \text{ is "small"} \end{array} \right\} \approx$$

$$\approx T \frac{y_{k+1} - y_k}{h}$$

Similarly, the force acting on (x_k, y_k) from the left point is approximately

$$T \frac{y_{k-1} - y_k}{h} \quad \text{so the total vertical force}$$

acting on x_k is

$$T \frac{y_{k+1} - 2y_k + y_{k-1}}{h}$$

So newtons law of acceleration gives

$$\frac{M}{L} h \frac{d^2 y_k(t)}{dt^2} = T \frac{y_{k+1} - 2y_k + y_{k-1}}{h}$$

That is if we write $y_{lc}(t) = y(x_k, t)$

$$\frac{\partial^2 y_k(t, x_k)}{\partial t^2} = \frac{IL}{M} \frac{y(t, x_k+h) - 2y(t, x_k) + y(t, x_k-h)}{h^2}$$

If we let $h \rightarrow 0$ we arrive at
(That is $N \rightarrow \infty$)

$$\frac{\partial^2 y(t, x)}{\partial t^2} = \frac{IL}{M} \frac{\partial^2 y(t, x)}{\partial x^2}, \quad \left. \right\} \text{Equations for a vibrating string}$$

Remarks:

1) Discuss this deduction 5 min.

2) Not mathematics!

i) too sloppy with Taylor expansions
and throwing away terms

ii) k is fixed but when $N \rightarrow \infty$
we "keep the same point" easy to fix.

iii) we assume that $\frac{y(t, x+h) - 2y(t, x) + y(t, x-h)}{h}$
converges

iv) We disregard forces in the x -direction
v) We do not formalize things mathematically
but rely on physics!

3) We will not attempt to do this with

mathematical stringency but assume that
 $\frac{\partial^2 y}{\partial t^2} = \lambda^2 \frac{\partial^2 y}{\partial x^2}$ is an important equation given by God.

Change of variables: Let

$$\frac{\partial^2 u(x,t)}{\partial t^2} = \lambda^2 \frac{\partial^2 u(x,t)}{\partial x^2} \quad \text{for all } x \in [0, L]$$

and $aX = x$, $bT = t$ and $U(X,T) = u(x,t)$

Then $U(X,T)$ is defined in $[0, \frac{L}{a}]$

and $\frac{\partial^2 U}{\partial T^2} = \left\{ \frac{\partial}{\partial T} \left[\frac{\partial t}{\partial T} \frac{\partial}{\partial t} + \frac{\partial X}{\partial T} \frac{\partial}{\partial X} \right] \right\} = \frac{1}{b^2} \frac{\partial^2 u(x,t)}{\partial t^2} =$

$$\frac{\lambda^2}{b^2} \frac{\partial^2 u(x,t)}{\partial x^2} = \left\{ \frac{\partial}{\partial x} \left[\frac{\partial X}{\partial x} \frac{\partial}{\partial X} + \frac{\partial T}{\partial x} \frac{\partial}{\partial T} \right] \right\} = \frac{\lambda^2 a^2}{b^2} \frac{\partial^2 U(X,T)}{\partial X^2}$$

so we can choose $a = \frac{L}{\pi}$ and $b = \frac{\pi}{\lambda L}$

and get a new equation of the same type

$$\frac{\partial^2 U(X,T)}{\partial T^2} = \frac{\partial^2 U(X,T)}{\partial X^2} \quad \text{in } [0, \pi]$$

It is therefore no loss of generality to assume that $L = \pi$ (or whatever else we fancy)

and $\lambda = 1$.

Important! In the future we will often, and without loss of generality, assume that a function $f(x)$ is defined on a

Question: Given some $f(x)$ and $g(x)$ defined on $[0, \pi]$

Can we find an $u(x,t)$ s.t.

$$\textcircled{1} \quad \frac{\partial^2 u(x,t)}{\partial t^2} = \frac{\partial^2 u(x,t)}{\partial x^2} \quad \text{in } [0, \pi] \times \{t > 0\}$$

$$\textcircled{2} \quad \partial u(x,0) = f(x) \quad \text{for } x \in [0, \pi]$$

$$\textcircled{3} \quad \frac{\partial u(x,0)}{\partial t} = g(x) \quad \text{for } x \in [0, \pi]$$

$$\textcircled{4} \quad u(0,t) = u(\pi,t) = 0 \quad \text{for } t > 0$$

Any ideas?

An absolutely crazy (and wonderfull) idea:

Assume that $u(x,t) = \varrho(x) \Psi(t)$

If $u(x,t) = \varrho(x) \Psi(t)$ then 1 Second,

$$\varrho(x) \Psi''(t) = \Psi(t) \varrho''(x) \Rightarrow \frac{\Psi''(t)}{\Psi(t)} = \frac{\varrho''(x)}{\varrho(x)} = -\lambda .$$

$$\Rightarrow \Psi''(t) = -\lambda \Psi(t) \Rightarrow \Psi(t) = A \cos(\sqrt{\lambda} t) + B \sin(\sqrt{\lambda} t)$$

$$\varrho''(x) = -\lambda \varrho(x) \quad \varrho(x) = C \cos(\sqrt{\lambda} x) + D \sin(\sqrt{\lambda} x)$$

Here we assume that $\lambda \geq 0$, a simple calculation shows that if $\lambda = -\mu^2 < 0$ then

$$\varrho(x) = A e^{\mu x} + B e^{-\mu x}$$

but using 4 gives

$$\varrho(0) = 0 \Rightarrow A = -B \quad \text{and}$$

$$\varrho(\pi) = A (e^{\mu \pi} - e^{-\mu \pi}) = 0 \Rightarrow A = 0 \quad \text{so } \varrho = 0$$

Now (4) states that $\varphi(0) = \varphi(\pi) = 0$

That is $A \cdot \cos(0) + B \cdot \sin(0) = A = 0$

and

$$\underbrace{A \cos(\sqrt{\lambda}\pi) + B \sin(\sqrt{\lambda}\pi)}_{=0} = 0$$

so if $B \neq 0$ then $\sqrt{\lambda} \in \mathbb{Z}$ so $\lambda = n^2$

for some $n \in \{0, 1, 2, \dots\}$.

So we may conclude that

$$u(x, t) = \sum_{n=0}^{\infty} B \sin(nx) (A_n \cos(nt) + B_n \sin(nt)). \quad (5)$$

But the equation is linear so we may add terms corresponding to different n and get a new solution

The solution in (5) satisfies (1) and (4)

But does it satisfy (2) and (3) ?

Condition ② states that

$$\sum_{n=0}^{\infty} A_n \sin(nx) = f(x) \quad (6)$$

and 3 that

$$\sum_{n=0}^{\infty} n B_n \sin(nx) = g(x) \quad (7)$$

Main question: (Fourier series)

Given a function $f(x)$ on $[0, \pi]$ can we write

$$f(x) = \sum_{n=0}^{\infty} A_n \sin(nx) + B_n \cos(nx) ?$$

So the absolutely crazy idea leads to whether we can write a function $f(x)$ as an infinite series of \sin & \cos .

This is an absolutely pregnant idea that has influenced very very much mathematics over the centuries.

Mathematics.

If $f(x) = \sum_{n=0}^{\infty} A_n \sin(nx) + B_n \cos(nx)$ for $x \in [0, \pi]$

then, given that the integrals exist,

$$\int_0^{\pi} f(x) \sin(kx) dx = \int_0^{\pi} \left(\sum_{n=0}^{\infty} A_n \sin(nx) \sin(kx) + \sum_{n=0}^{\infty} B_n \cos(nx) \sin(kx) \right) dx =$$

$$= A_k \cdot \frac{\pi}{2} \quad \text{and similarly} \Rightarrow A_k = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(kx) dx$$

$$\boxed{B_k = \frac{2}{\pi} \int_0^{\pi} f(x) \cos(kx) dx}$$

Questions: C

1) For what functions (if any) does

$$f(x) = \lim_{N \rightarrow \infty} \sum_{n=0}^N (A_n \cos(nx) + B_n \sin(nx))$$

and in what sense?

Where

$$A_n = \frac{2}{\pi} \int_0^{\pi} f(x) \cos(nx) dx$$

$$B_n = \frac{2}{\pi} \int_0^{\pi} f(x) \sin(nx) dx$$

Fourier analysis

Q:
How far
what can
be said of
the set where
 $\sum A_n \sin nx + B_n \cos nx$
is disjoint

Center

Transfinite
sets

View
 $\sin(nx)$ & $\cos(nx)$
as a basis
for a vector
space

Hilbert

Functional
Analysis

General
Convergence
questions

Analytic
Riemann integral

What
do
differential
operators

Analytic
number
theory