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## $\infty$ -Ground states in the plane

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based on joint work with Peter Lindqvist

### Archipelagic perspectives on mathematics, physics

and perceptible spectra of reality

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Introduction Introduction to Δ∞ Eigenvalue problem: from finite p to ∞ Results Ideas and tools Open problems What will I talk about?

The  $\infty$ -eigenvalue equation:

$$\max\left\{\lambda - \frac{|\nabla u|}{u}, \underbrace{\sum_{i,j} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \frac{\partial^2 u}{\partial x_i \partial x_j}}_{\Delta_{\infty} u}\right\} = 0$$

Arises as the Euler-Lagrange equation of the Rayleigh quotient

$$\frac{\|\nabla u\|_{L^{\infty}(\Omega)}}{\|u\|_{L^{\infty}(\Omega)}}$$



The problem is a highly nonlinear version of the eigenvalue problem for the Laplacian:

Minimizers of

$$\frac{\int_{\Omega} |\nabla u|^2}{\int_{\Omega} |u|^2}, \quad \text{with } u = 0 \text{ on } \partial\Omega$$

satisfy

$$\Delta u + \lambda u = \mathbf{0}$$

where  $\lambda$  is the minimum.

 $\Delta_\infty: The infinity Laplacian$ 

The infinity Laplacian

$$\Delta_{\infty} u := \langle \nabla u, D^2 u \, \nabla u \rangle = \sum_{i,j=1}^{n} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \frac{\partial^2 u}{\partial x_i \partial x_j}$$

Solutions of

$$\Delta_{\infty} u = 0$$

are called  $\infty$ -harmonic functions.

Discovered by Gunnar Aronsson in the 60's in connection to Lipschitz extensions.

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If *u* minimizes

$$\int_{\Omega} |\nabla u|^2, \text{ among functions coinciding on } \partial \Omega,$$

then u is harmonic and

 $\Delta u = 0$  in  $\Omega$ .

 $\Delta_{p} - The p-Laplacian$ 

If  $u_p$  ( $p \ge 2$ ) minimizes

$$\int_{\Omega} |\nabla u|^{p}, \text{ among functions coinciding on } \partial\Omega,$$

then  $u_p$  is *p*-harmonic and

$$\Delta_{
ho} u_{
ho} = {
m div}(|
abla u_{
ho}|^{
ho-2}
abla u_{
ho}) = 0 \quad {
m in} \; \Omega$$

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# $\Delta_{\infty}$ via the *p*-Laplacian

As  $p \to \infty$  $\|\nabla u\|_{L^p(\Omega)} \to \|\nabla u\|_{L^\infty(\Omega)},$  $\Delta_p u = |\nabla u|^{p-2} \Delta u + (p-2)|\nabla u|^{p-4} \Delta_\infty u \to \Delta_\infty u$ 

#### Reasonable that $u_p \rightarrow u$ where u minimizes

 $\|\nabla u\|_{L^{\infty}(\Omega)}$ , among functions coinciding on  $\partial \Omega$ 

and solves  $\Delta_{\infty} u = 0$  in  $\Omega$ .

Aronsson 66. Bhattacharya, DiBenedetto and Manfredi 89.

$$\begin{array}{c} & \\ \text{Introduction} \\ \text{Introduction to } \Delta_{\infty} \\ \\ \text{Eigenvalue problem: from finite } p \text{ to } \infty \\ \\ \text{Results} \\ \text{Ideas and tools} \\ \\ \text{Open problems} \end{array}$$

Let  $\Omega$  be open and bounded,  $g:\partial\Omega 
ightarrow \mathbb{R}$  be Lipschitz and

$$egin{cases} \Delta_\infty u = 0 & ext{in } \Omega \ u = g & ext{on } \partial \Omega. \end{cases}$$

Then

$$\sup_{x,y\in\Omega}\frac{|u(x)-u(y)|}{|x-y|} = \sup_{x,y\in\partial\Omega}\frac{|g(x)-g(y)|}{|x-y|}$$

and

$$\|\nabla u\|_{L^{\infty}(\Omega)} \leq \|\nabla g\|_{L^{\infty}(\partial\Omega)}$$

This was first proved by Aronsson for  $C^2$  functions.



- Classical solutions is not a good notion of solutions. One should use viscosity solutions.
- Existence and uniqueness of solutions of the Dirichlet problem on bounded domains, Aronsson 67, Jensen 93.
- Solutions vs minimizers of ||∇u||, Aronsson 67, Crandall-Evans-Gariepy 2001
- Differentiability in any dimension, Evans-Smart 2011
- $C^{1,\alpha}$ -regularity in the plane, Savin-Evans 2008

• 
$$C^2 + u_{xx}u_{yy} - u_{xy}^2 \neq 0 \Rightarrow C^{\infty}$$
, Aronsson 67.

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# Some $\infty$ -harmonic functions

• Cones:  $|x - x_0|$  for  $x \neq x_0$ 

- Aronsson's function  $x^{\frac{4}{3}} y^{\frac{4}{3}}$ . It is merely  $C^{1,1/3}$  which is believed to be the optimal regularity of solutions.
- Any  $C^1$  solution of the eikonal equation  $|\nabla u| = \text{constant}$ . Note that

$$\frac{1}{2}\Delta_{\infty}u = \langle \nabla u, \nabla |\nabla u|^2 \rangle.$$

• The distance function to a set is  $\infty$ -harmonic wherever it is  $C^1$ .



The functional for  $\Delta_{\infty}$  is not additive (as the one for  $\Delta$ ) so we ask that it is a minimizer on any subdomain, otherwise the set of minimizers can be large.

**Example:** A stadium minus a point with boundary data identically equal to one on the boundary of the stadium and 0 at the removed point.



Friedrichs's inequality (for p > 1 and  $\Omega$  bounded):

 $\|u\|_{L^p(\Omega)} \leq C \|\nabla u\|_{L^p(\Omega)}$ 

for smooth functions vanishing on  $\partial \Omega$ .

The associated Rayleigh quotient:

$$\lambda_{p} = \inf_{u \in W_{0}^{1,p}(\Omega)} \frac{\int_{\Omega} |\nabla u|^{p}}{\int_{\Omega} |u|^{p}}$$

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## The eigenvalue equation for finite *p*

# The eigenvalue equation:

Minimizers of

$$\frac{\int_{\Omega} |\nabla u|^{p}}{\int_{\Omega} |u|^{p}}, \quad u \in W_{0}^{1,p}(\Omega)$$

satisfy

$$\Delta_p u + \lambda_p |u|^{p-2} u = 0$$

**Terminology:** A *ground state* is a minimizer of the Rayleigh quotient.



- Ground states equivalent to solutions.
- The ground state is unique up to a multiplicative constant. Thelin (balls), Sakaguchi (convex domains), Anane (C<sup>2,α</sup>-domains), Lindqvist (any).
- The ground state is log-concave (convex domains). (Sakaguchi generalized the Brascamp-Lieb Theorem)
- The first eigenvalue is isolated and there is a well-defined second eigenvalue.
- **Unknown** if the eigenvalues are countable ( $p \neq 2$ ).

**Proof for** p = 2: (Belloni-Kawohl, 2002, general p): Consider two extremals with u, v with

$$||u||_{L^2(\Omega)} = ||v||_{L^2(\Omega)} = 1.$$

Put

$$w = \left(\frac{u^2 + v^2}{2}\right)^{\frac{1}{2}}$$

Then

$$\int_{\Omega} w^2 dx = 1.$$

By the convexity of  $x \mapsto x^2$ 

$$\begin{split} \int |\nabla w|^2 &= \frac{1}{2} \int_{\Omega} (u^2 + v^2) \left( \frac{u^2}{u^2 + v^2} \frac{\nabla u}{u} + \frac{v^2}{u^2 + v^2} \frac{\nabla v}{v} \right)^2 \\ &\leq \frac{1}{2} \int_{\Omega} |\nabla u|^2 + |\nabla v|^2, \end{split}$$

so that

$$\lambda_2 \leq \int_{\Omega} |\nabla w|^2 \leq \frac{1}{2} \int_{\Omega} |\nabla u|^2 + |\nabla v|^2 \leq \lambda_2.$$

The strict convexity of  $x \mapsto x^2$  forces w = u = v.

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The  $\infty$ -eigenvalue problem

The Rayleigh quotient

$$\left(\inf_{u\in W_0^{1,\rho}(\Omega)}\frac{\int_{\Omega}|\nabla u|^{\rho}}{\int_{\Omega}|u|^{\rho}}\right)^{\frac{1}{\rho}}\to \inf_{u}\frac{\|\nabla u\|_{L^{\infty}(\Omega)}}{\|u\|_{L^{\infty}(\Omega)}}:=\lambda_{\infty}$$

The eigenvalue equation:

$$\max\left\{\lambda_{\infty} - \frac{|\nabla u|}{u}, \underbrace{\sum_{i,j} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} \frac{\partial^2 u}{\partial x_i \partial x_j}}_{\Delta_{\infty} u}\right\} = 0$$

First studied by Juutinen, Lindqvist, Manfredi.



- $\lambda_p^{\frac{1}{p}} \to 1/R := \lambda_{\infty}$ , where *R* is the radius of the largest ball that can be inscribed in  $\Omega$ .
- There are many more minimizers than solutions.
- The first egenfunction is unique in a large class of domains including the ball and stadiums, Yu. Unknown in general.
- There is a counter example to uniqueness in a non-convex, dumbbell shaped domain with at least three linearly independent ground states. Hynd-Smart-Yu.

We call a non-negative solution of the equation (with  $\lambda = 1/R$ ) a ground state.



- In a ball (or a stadium) the ground state is the distance function.
- In the square, the distance function is *not* a ground state, however

$$v \le u \le d$$

where v is the  $\infty$ -harmonic function which is 1 at the midpoint and zero on the boundary.



To state our results, we assume

- $\Omega \subset \mathbb{R}^2$  is a convex polygon
- *u* is a limit of *p*-eigenvalues (*variational* ground state)

Then (Yu and Sakaguchi):

- In u is concave
- u is  $\infty$ -harmonic outside a closed set  $\Upsilon$  with zero measure
- the set of singular points of u is contained in Υ
- $|\nabla u| > \lambda_{\infty} u$  outside  $\Upsilon$
- on  $\Upsilon$  we have " $|\nabla u| = \lambda_{\infty} u$ "
- *u* attains its max exactly on the set where the distance function attains its max (the high ridge)



A streamline  $\alpha = \alpha(t)$ , is a solution of

$$\frac{d\alpha}{dt} = \nabla u(\alpha(t)).$$

If u is  $\infty$ -harmonic

$$\frac{d}{dt}|\nabla u(\alpha(t))|^2 = 2\Delta_{\infty}u(\alpha(t)) = 0.$$

Hence,

$$|\nabla u(\alpha(t))| =$$
 the speed = constant.

Requires second order derivatives! In general not true.



Let *H* be the *high ridge*. Streamlines starting at a corner are called *attracting* streamlines and streamlines starting at the point where  $|\nabla u|$  attains a max between two corners is called a *median*. The convexity implies that the normal derivative  $(= |\nabla u|)$  is monotone along the half-edges corner-median.

**Theorem 1:**  $\Upsilon$  lies in set of attracting streamlines. Streamlines starting at other points cannot meet before joining an attracting streamlines or *H*. Along such streamlines, the speed  $|\nabla u|$  is constant.

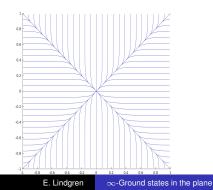


**Cor 1:** The arc of a streamline from the boundary to an attracting streamline (or *H*) is either convex or concave. The length of this arc is  $u/|\nabla u|$  at the hitting point. In particular, all arcs with the first meeting point in  $\Upsilon$  have unit length.

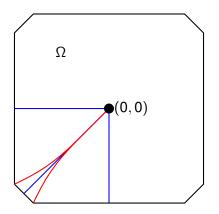
**Cor 2:** A median is a straight line segment until it joins some attracting streamline (or *H*).



Let  $\Omega$  be a square and K the center. The attracting streamlines are the four half-diagonals. All streamlines meet at a diagonal, except the four segments along the coordinate axes.



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The speed along a streamline  $\alpha$ 

$$\left|\frac{d\alpha(t)}{dt}\right| = |\nabla u_{\infty}(\alpha(t))|$$

is non-decreasing outside  $\Upsilon$ .



Assume that  $D \subset \subset \Omega \setminus \Upsilon$  and p > 2. Then

$$\oint_{\partial D} |\nabla u|^{p-2} \langle \nabla u, \mathbf{n} \rangle \, ds \, \leq \, 0$$

where **n** is the outer normal.

**Idea:** If for  $u_p$ , then the inequality follows from that

$$\oint_{\partial D} |
abla u_{
ho}|^{p-2} \langle 
abla u, \mathbf{n} 
angle \, ds = \int_{D} \Delta_{
ho} u_{
ho} \, dx \leq 0,$$

since  $\Delta_{\rho}u_{\rho} \leq 0$ .

**Consequence:** The gradient is non-increasing along streamlines as long as they do not meet.



## Sketch of proof:

Assume that

- the points  $x_1$  and  $x_2$  are on the same level curve u = a,
- the points  $y_1$  and  $y_2$  both are on the higher level curve u = b > a,
- ascending streamlines join x<sub>1</sub> with y<sub>1</sub> and x<sub>2</sub> with y<sub>2</sub> but do not meet before.

Then

$$\|\nabla u\|_{\infty,\overline{y_1y_2}} \leq \|\nabla u\|_{\infty,\overline{x_1x_2}},$$

that is, the lower level curve has the larger gradient.



**The idea** is to exploit that the speed is non-decreasing and that sometimes, when the streamlines do not meet, it is also non-increasing, so that it is constant along suitable arcs of streamlines.

Note also that since  $\log u$  is concave,  $|\nabla u|/u$  is decreasing along streamlines. So once a streamline hits the set  $\Upsilon$ , the rest of the streamline also lies in  $\Upsilon$ .



- When is a ground state ∞-harmonic also across the attracting streamlines?
- Can we prove uniqueness using this in some simple cases?
- General (smooth) convex sets?
- We only use *log*-concavity, not that we have a limit of  $u_p s$ .

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## Thank you for listening!



#### Some references:

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